

$$\gamma_g \log(2) = \lambda_g \log(2) + \nu_2(2i\pi)$$

The shape of the interaction region of colliding protons in a Regge model

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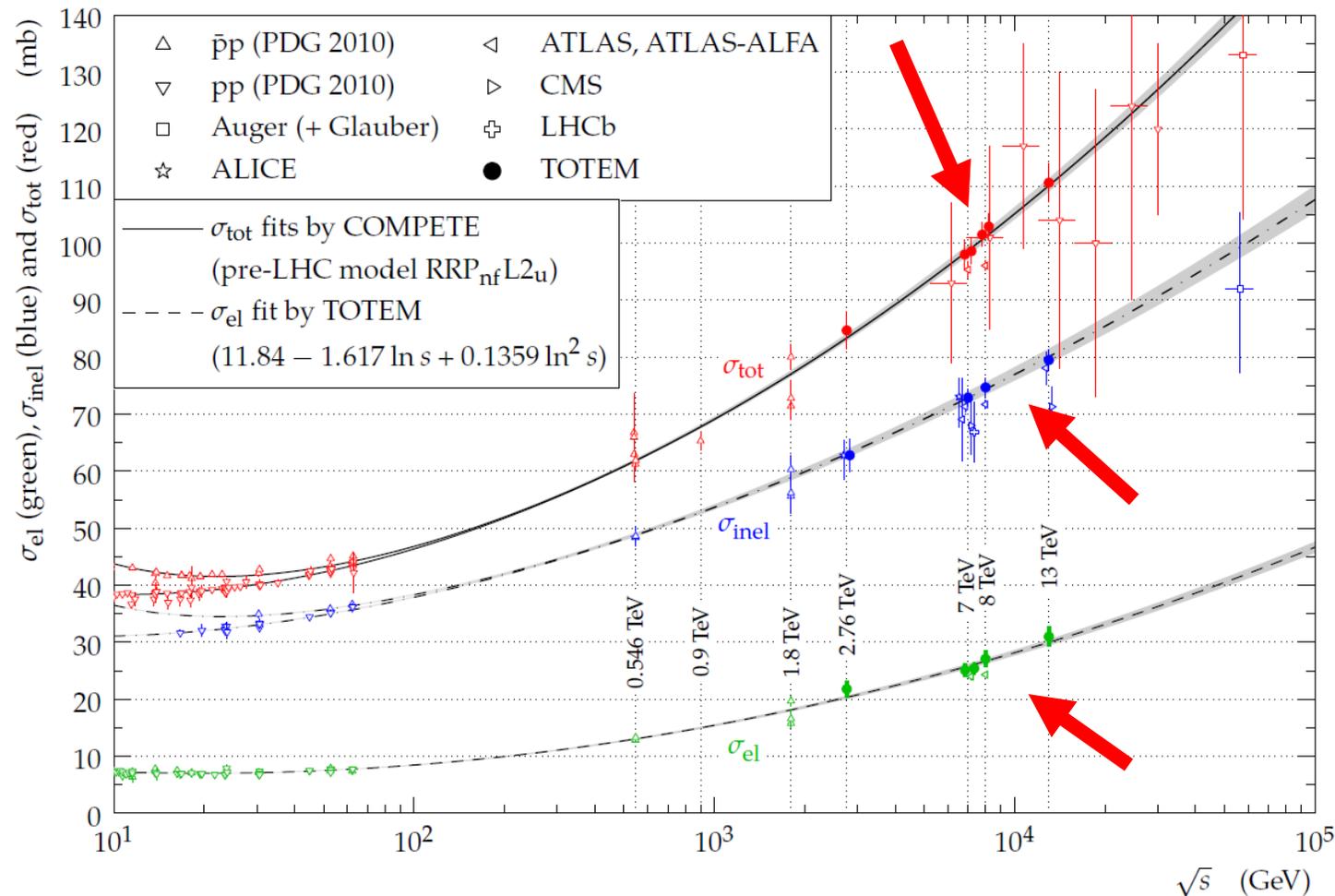
New Trends in High-Energy Physics Conference

May 12 – 18, 2019, Odessa, Ukraine



Introduction

- Total, elastic and inelastic cross sections follow the expectations from lower energies.



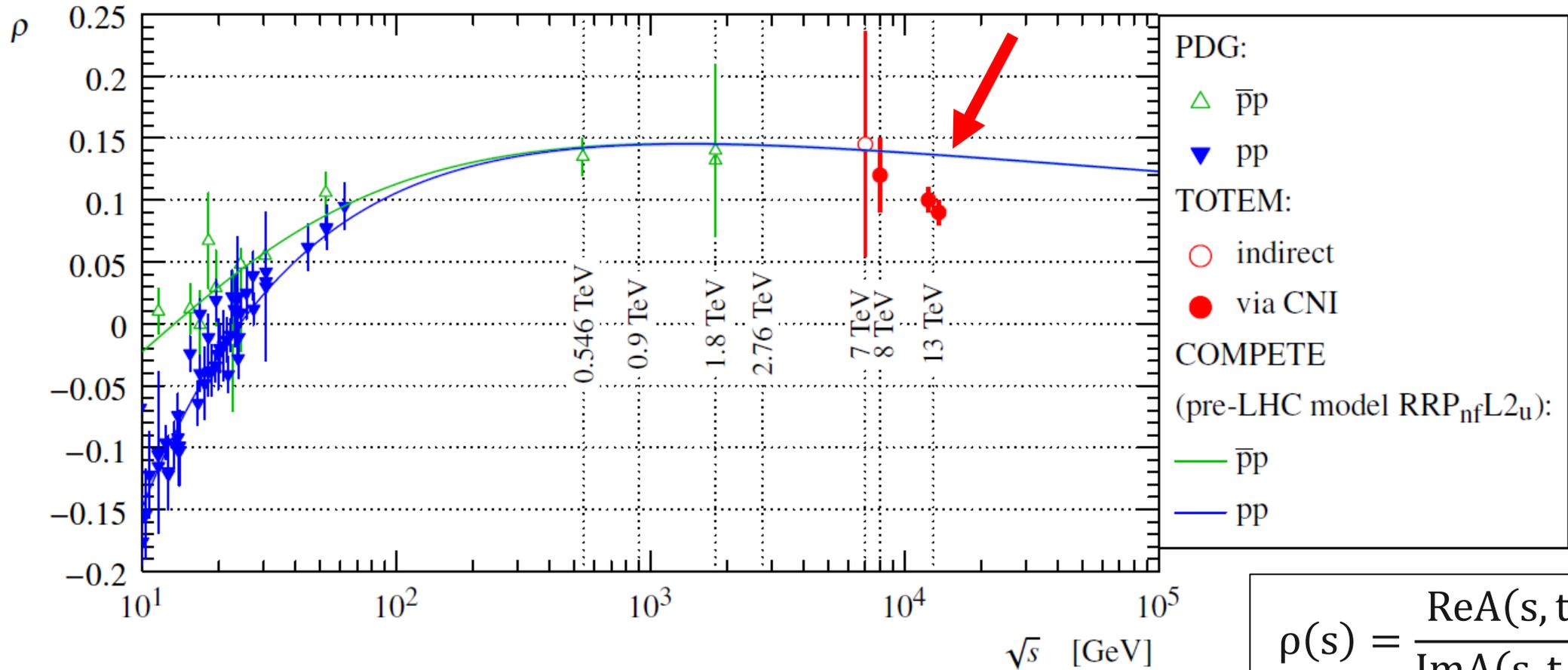
[TOTEM Collab.,
Eur. Phys. J. C 79
\(2019\) 103](#)

pp and $\bar{p}p$ total, elastic and inelastic cross sections measurements.

Introduction

- Very low value of the pp ρ -parameter at 13 TeV.

[TOTEM Collab., CERN-EP-2017-335](#)

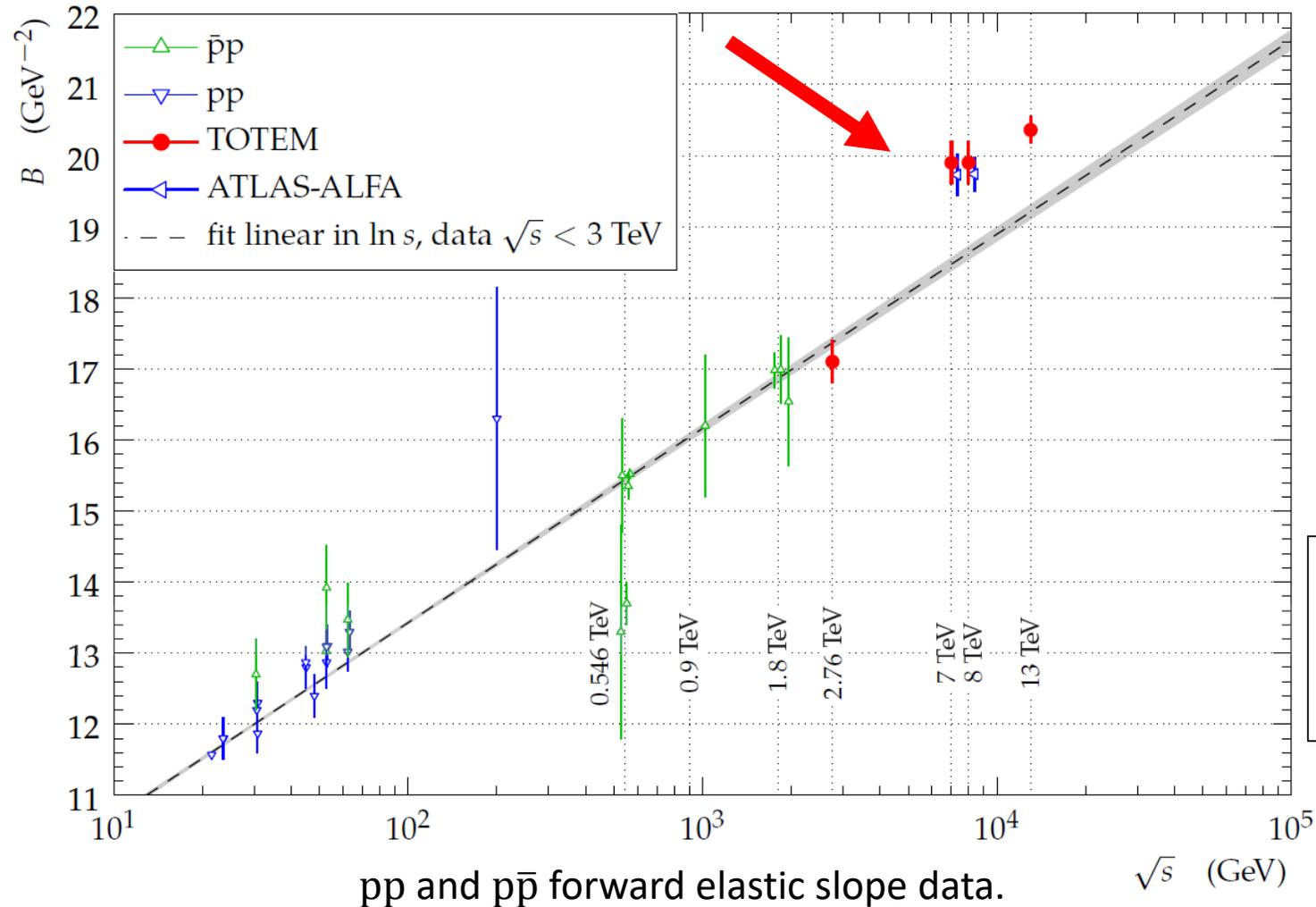


pp and $\bar{p}p$ ρ -parameter measurements.

$$\rho(s) = \frac{\text{Re}A(s, t = 0)}{\text{Im}A(s, t = 0)}$$

Introduction

- Rapid rise of the pp forward slope from about 3 TeV.



[TOTEM Collab.,
Eur. Phys. J. C 79 \(2019\) 103](#)

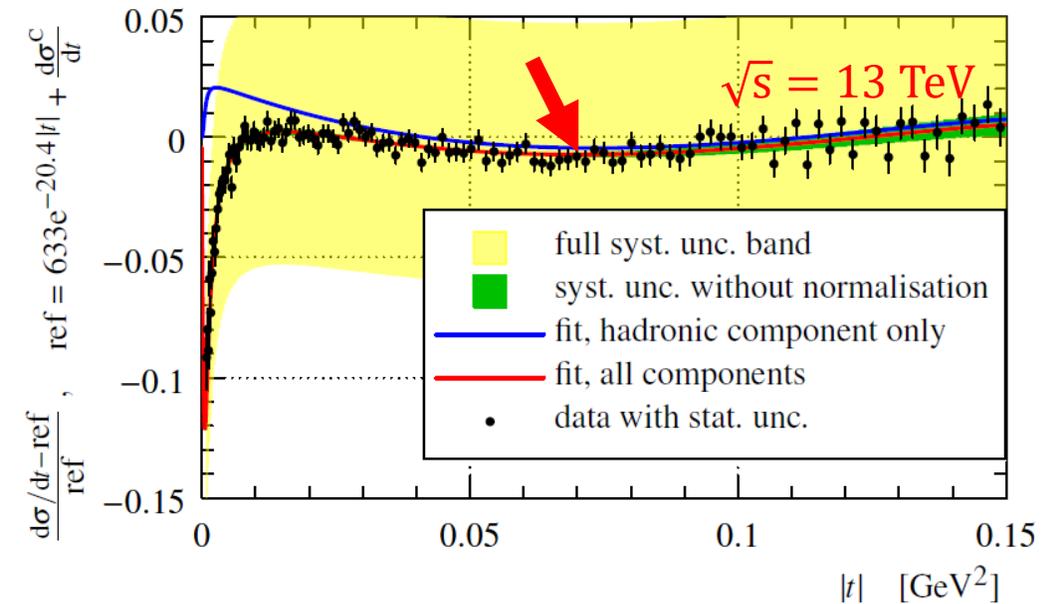
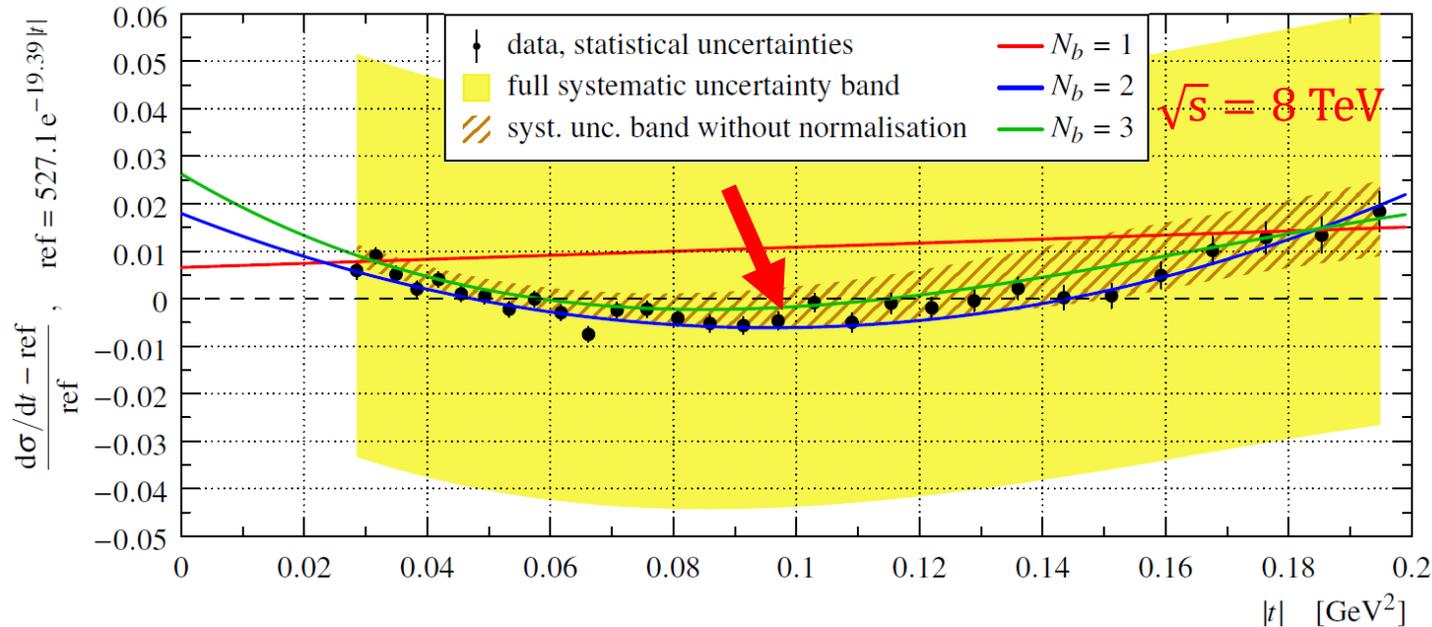
$$B(s) = \left. \frac{d}{dt} \left(\ln \frac{d\sigma_{el}}{dt}(s, t) \right) \right|_{t=0}$$

Introduction

- Low- $|t|$ structure in the pp differential cross section (the so called "break").

[TOTEM Collab., Nucl. Phys. B 899 \(2015\) 527-546](#)

[TOTEM Collab., CERN-EP-2017-335](#)



Normalized form of the elastic pp low- $|t|$ differential cross section data measured by TOTEM at 8 and 13 TeV.

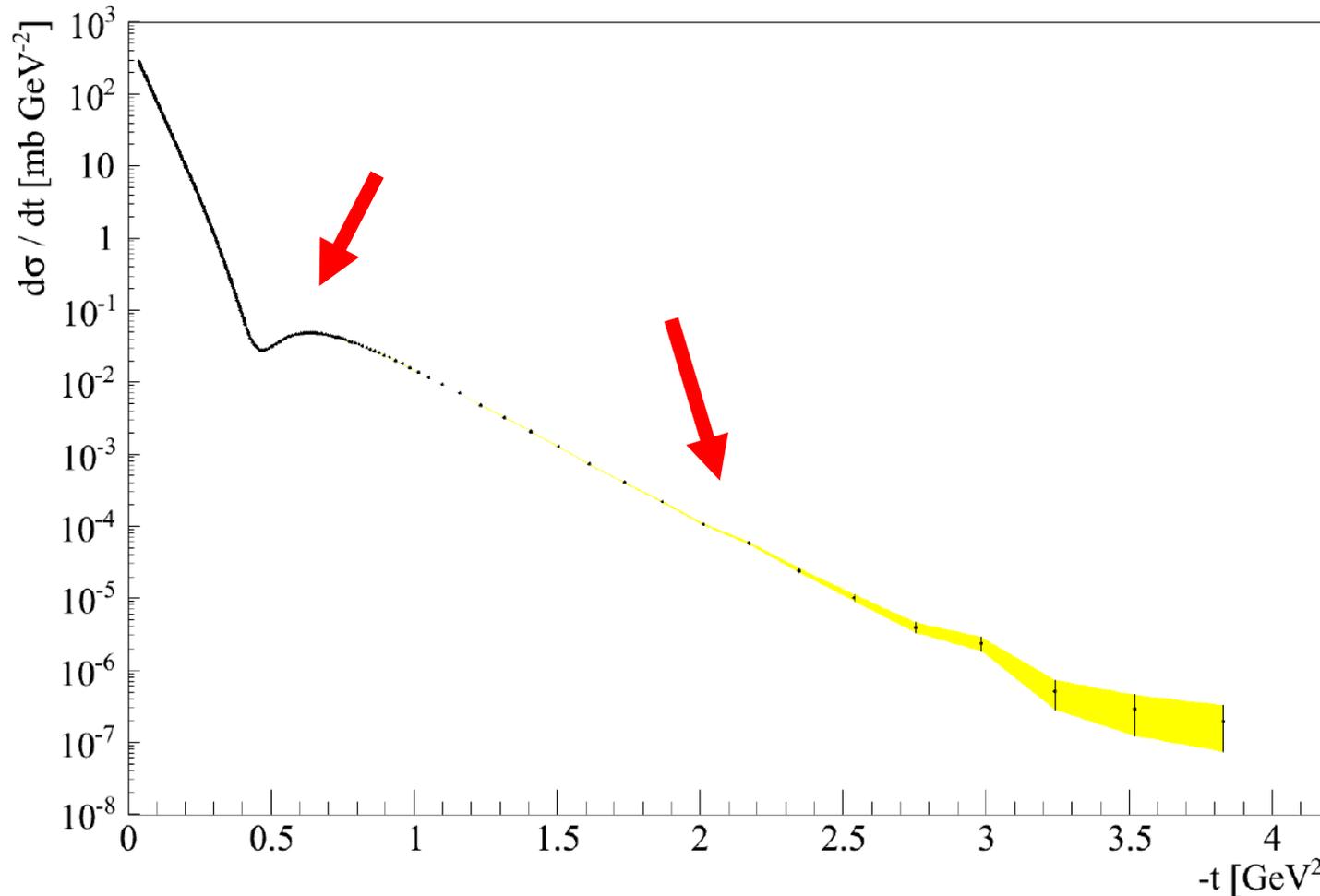
L. Jenkovszky, I. Szanyi, Phys. Part. Nuclei Lett., 14 (2017) 687

L. Jenkovszky, I. Szanyi, Mod. Phys. Lett. A 32, 1750116 (2017)

L. Jenkovszky, I. Szanyi, Cl. Tan, Eur. Phys. J. A 54: 116 (2018)

Introduction

- Absence of secondary dips and bumps in the pp differential cross section.



TOTEM Collab.,
CERN-EP-2018-338

Elastic pp differential cross section preliminary data measured by TOTEM at 13 TeV.

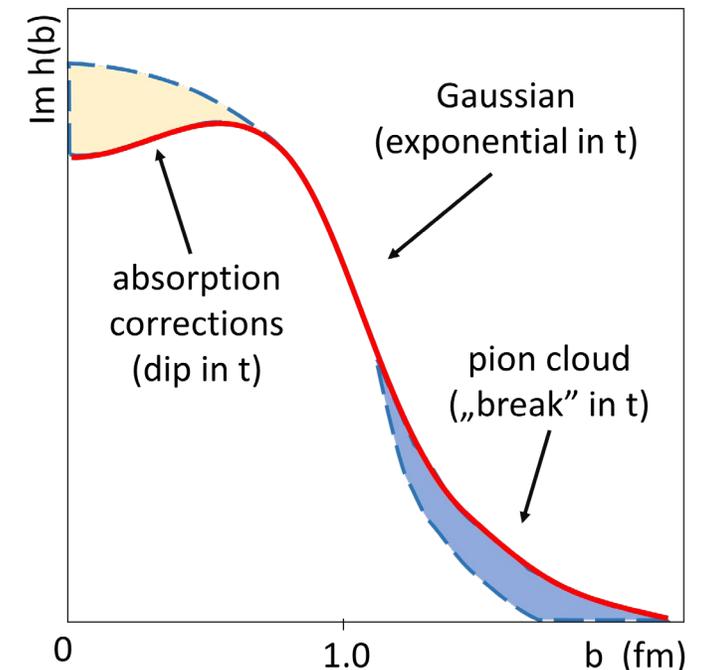
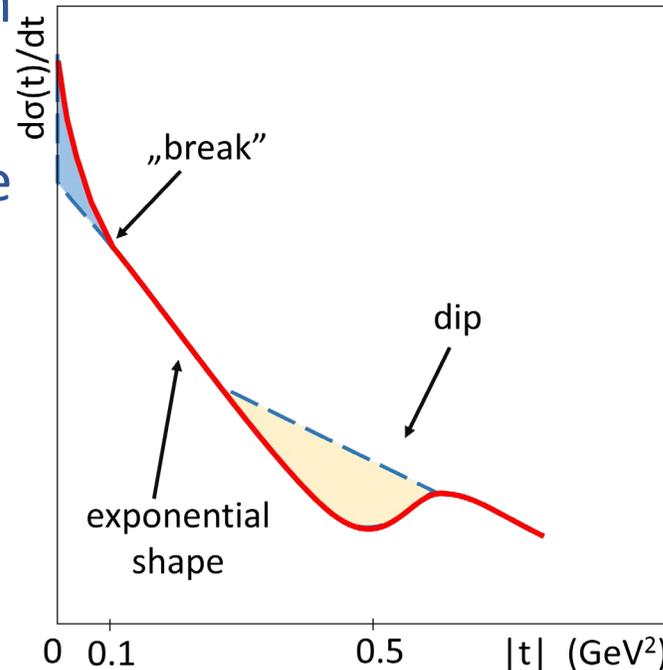
The $d\sigma(t)/dt$ and the $\text{Im } h(b)$

$$\frac{d\sigma_{\text{el}}}{dt}(s, t) = \frac{\pi}{s^2} |A(s, t)|^2$$



$$h(s, b) = \frac{1}{s} \int_0^{\infty} A(s, t) J_0(b\sqrt{-t}) \sqrt{-t} dt$$

- **„break“**: deviation from an exponential form near $-t = 0.1 \text{ GeV}^2$
 - related to the two-pion exchange (t-channel unitarity)
 - pion cloud of the proton
- **dip**: diffraction minimum with energy dependent location
 - related to absorption corrections (s-channel unitarity)
 - decrease in $\text{Im } h(b)$ at small b



Connection between the elastic pp differential cross section and the imaginary part of the impact parameter amplitude (schematic view).

Scattering amplitude: dipole Regge model

$$A(s, t)_{\bar{p}p} = A_P(s, t) \pm A_O(s, t) + A_f(s, t) \pm A_\omega(s, t) + \dots$$

- Dipole pomeron and odderon:**

$$A_P(s, t) = i \frac{a_P s}{b_P s_{0P}} [r_{1P}^2(s) e^{r_{1P}^2(s)[\alpha_P - 1]} - \epsilon_P r_{2P}^2(s) e^{r_{2P}^2(s)[\alpha_P - 1]}]$$

[I. Szanyi, N. Bence, L. Jenkovszky: J. Phys. G 46, 055002 \(2019\)](#)

$$A_O(s, t) = -i A_P(s, t)$$

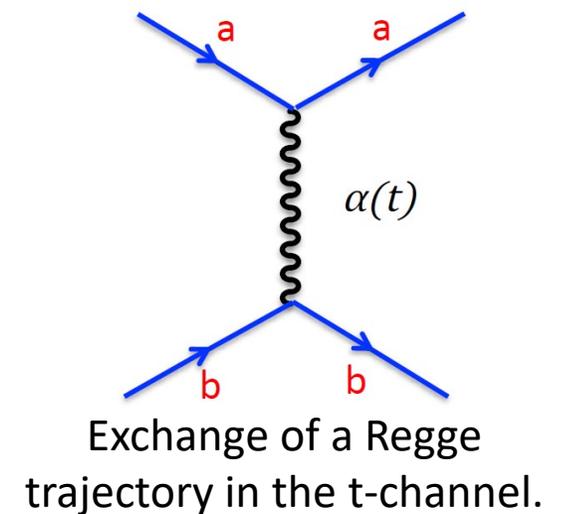
(with free parameters labeled by "O")

$$r_{1P}^2(s) = b_P + L_P - i\pi/2 \quad r_{2P}^2(s) = L_P - i\pi/2 \quad L_P \equiv \ln(s/s_{0P})$$

- Pomeron and odderon trajectories:**

$$\alpha_P \equiv \alpha_P(t) = 1 + \delta_P + \alpha_{1P}t - \alpha_{2P} \left(\sqrt{4m_\pi^2 - t} - 2m_\pi \right)$$

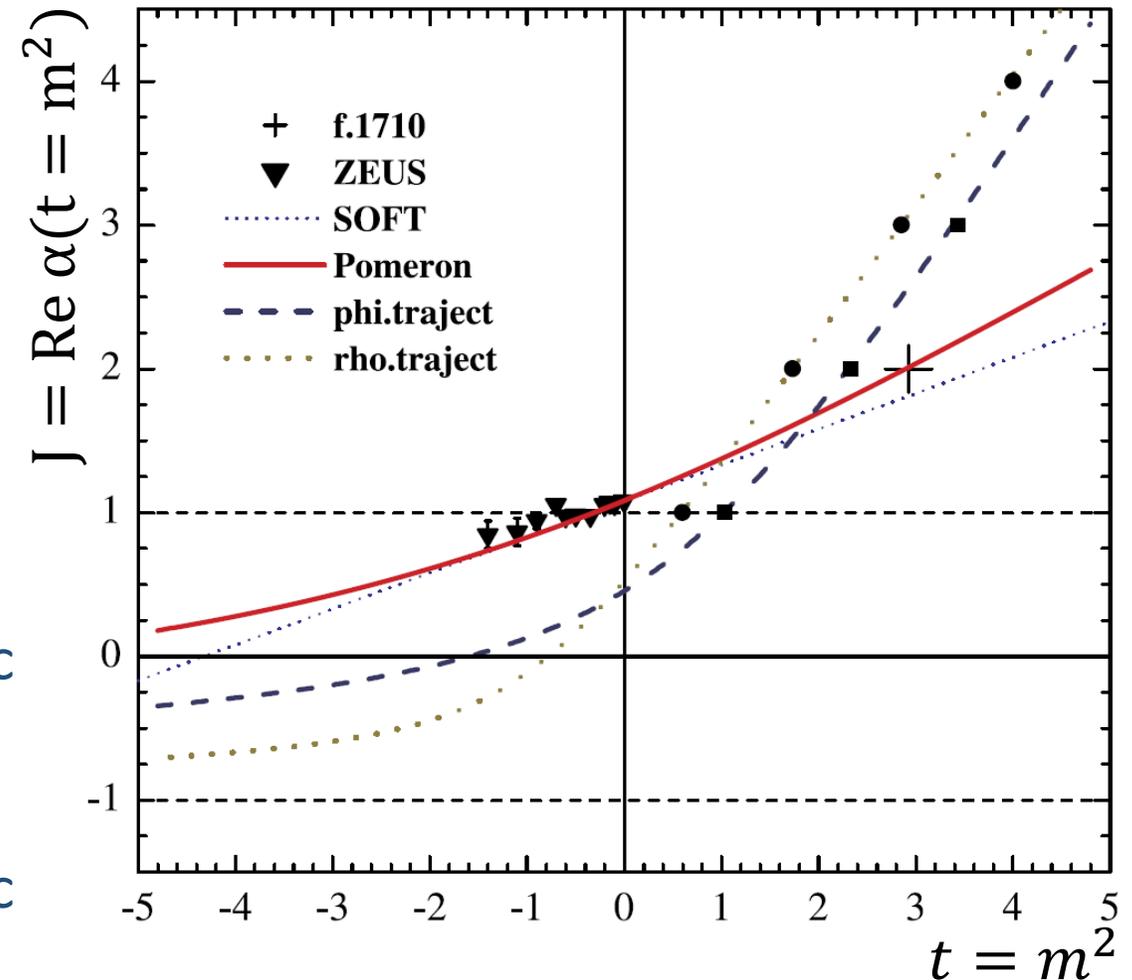
$$\alpha_O \equiv \alpha_O(t) = 1 + \delta_O + \alpha_{1O}t - \alpha_{2O} \left(\sqrt{9m_\pi^2 - t} - 3m_\pi \right)$$



Regge trajectories

- Reggeon (in general):
 - virtual particle with continuously varying spin ($J = \text{Re } \alpha(t)$) and virtuality ($t = m^2$) lying on the relevant trajectory (scattering at $-t$)
 - at certain values of virtuality there are real particles (spectroscopy at $+t$)
- Secondary reggeons ($f, \omega, \rho, \phi \dots$) → mesonic exchanges and meson spectra
- Pomeron (P) and odderon (O) → gluonic exchanges and glueball spectra

L. Jenkovszky, R. Schicker and I. Szanyi.
Int. J. Mod. Phys. E 27, 1830005 (2018)



The Chew–Frautschi plots of the leading ρ , ϕ and P complex Regge trajectories
 M.N. Sergeenko, Eur. Phys. J. C (2012) 72:2128

Unitarity and overlap functions

- **Unitarity of the S-matrix**, $SS^+ = 1 \rightarrow$ relation between the elastic scattering amplitude and the inelastic processes.
- **The unitarity constraint in impact parameter representation:**

$$2\text{Im}h(s, b) = |h(s, b)|^2 + G_{\text{in}}(s, b)$$

- **The total, elastic and inelastic cross sections:**

$$\sigma_{\text{tot}}(s) = 2 \int d^2b \text{Im}h(s, b)$$

$$\sigma_{\text{el}}(s) = \int d^2b |h(s, b)|^2$$

$$\sigma_{\text{in}}(s) = \int d^2b G_{\text{in}}(s, b)$$

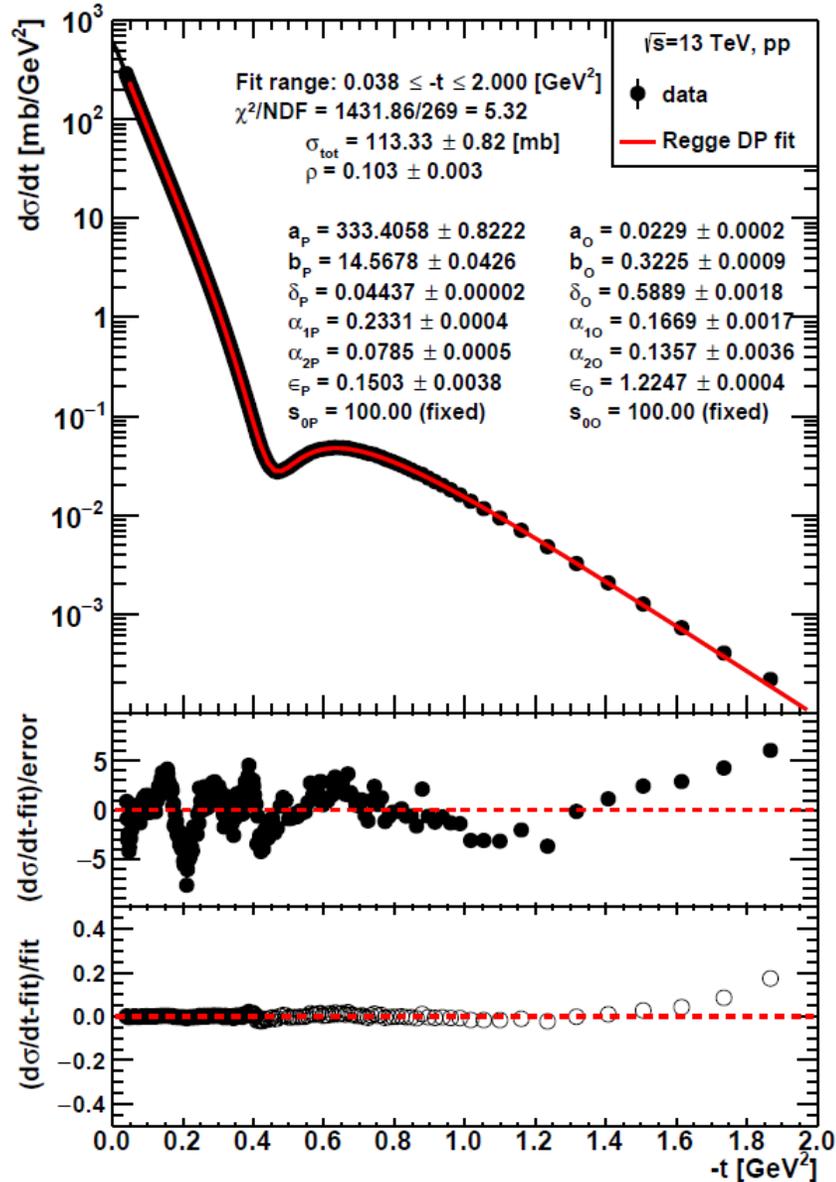
- **The inelastic overlap function G_{in}** \rightarrow probability of absorption associated to a given b value \rightarrow gives the shape of the interaction region:

$$G_{\text{in}}(s, b) = 2\text{Im}h(s, b) - |h(s, b)|^2$$

with

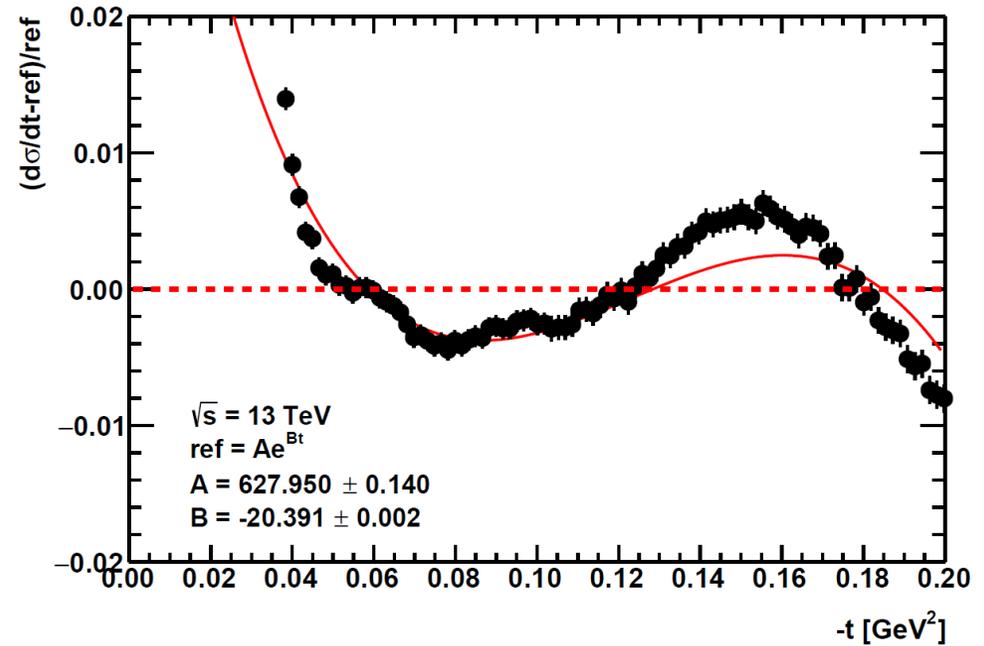
$$0 \leq G_{\text{in}}(s, b) \leq 1$$

Fit of $d\sigma(t)/dt$ @ 13 TeV



$$R(t) = \frac{d\sigma/dt - \text{ref}}{\text{ref}}$$

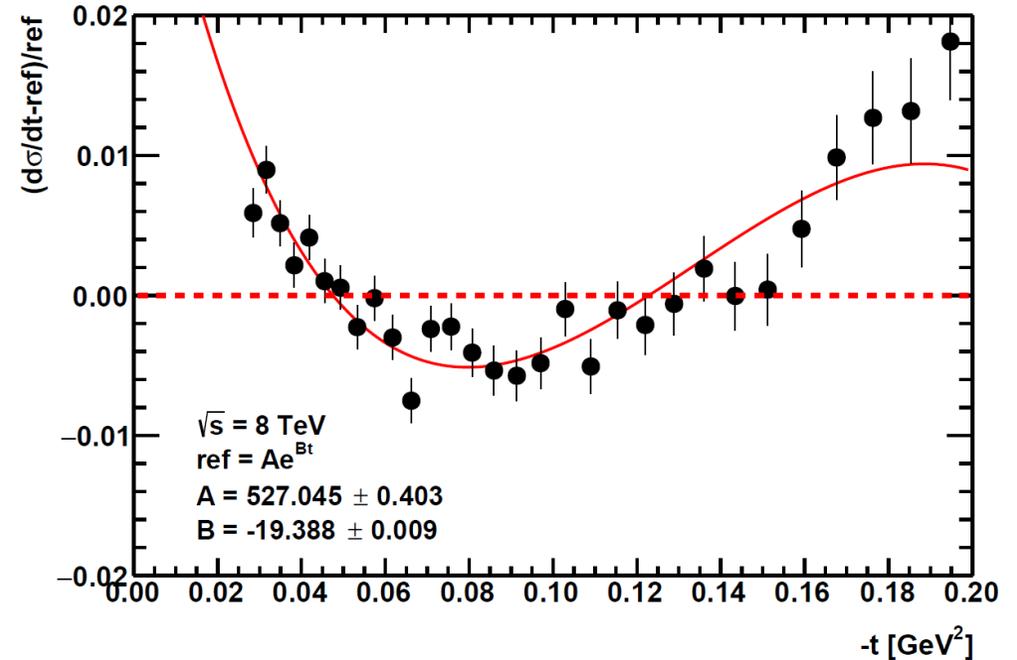
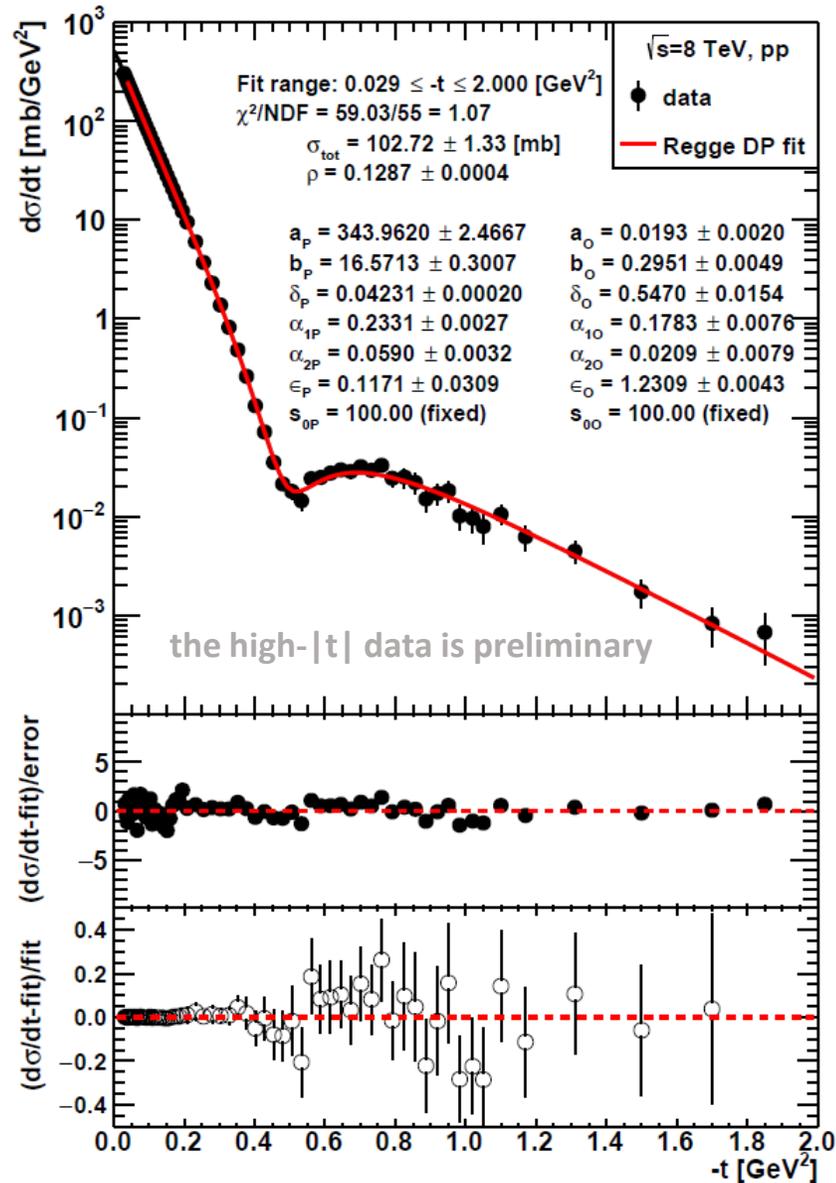
$$\text{ref} = Ae^{Bt}$$



Normalized form of the differential cross section showing the low- $|t|$ "break"

Source of the data:
TOTEM Collab., CERN-EP-2018-338

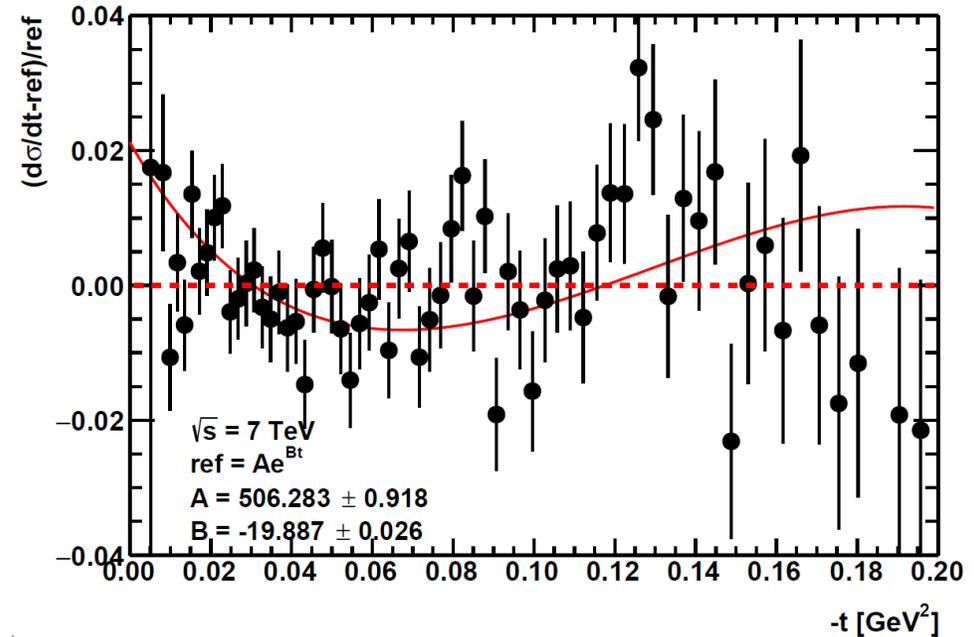
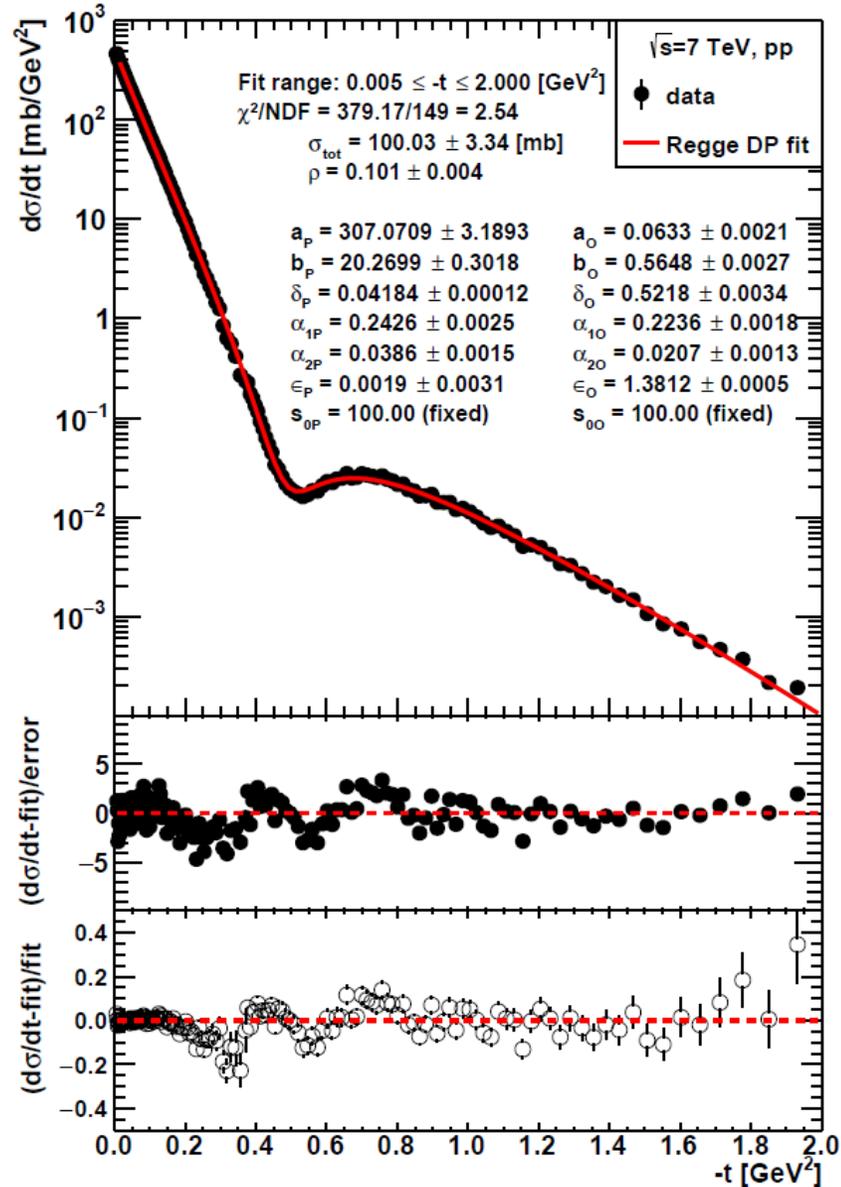
Fit of $d\sigma(t)/dt$ @ 8 TeV



Normalized form of the differential cross section showing the low- $|t|$ "break"

Source of the data:
 TOTEM Collab., Nucl. Phys. B 899 (2015) 527-546
 J. Kaspar, ISMD2017, Tlaxcala, Mexico (2017).

Fit of $d\sigma(t)/dt$ @ 7 TeV

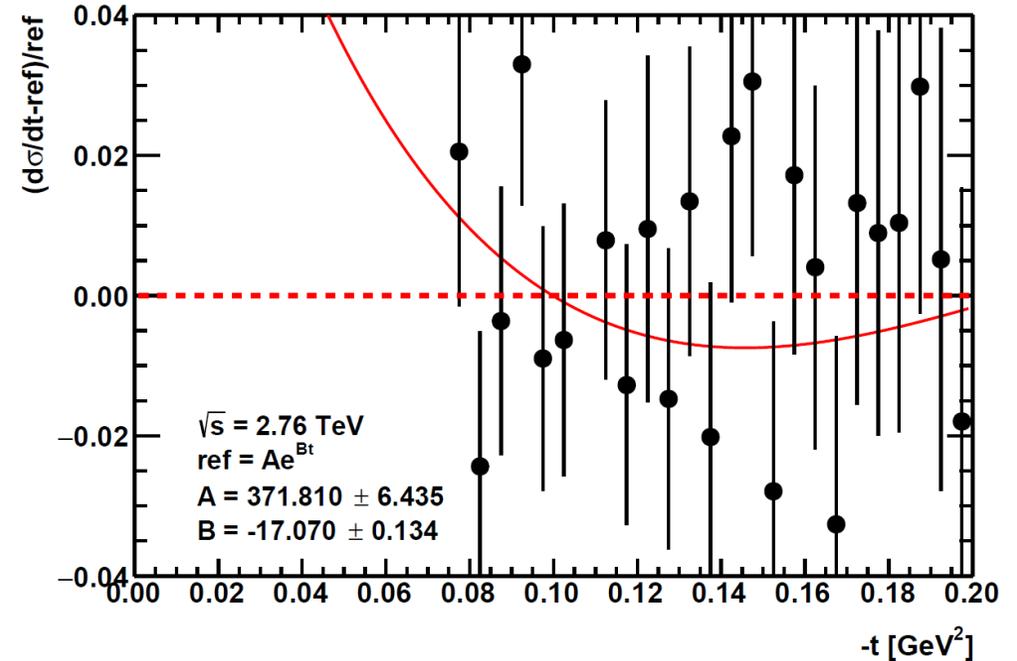
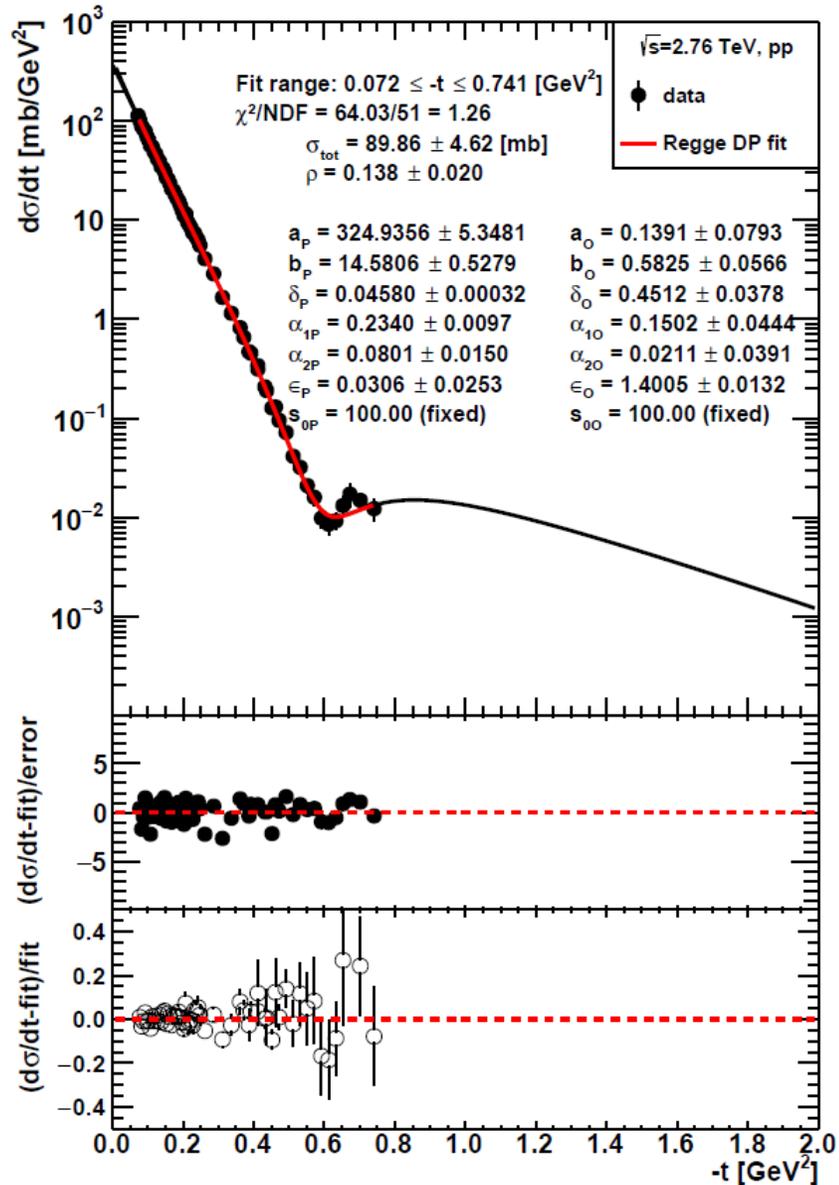


Normalized form of the differential cross section showing the low- $|t|$ "break"

Source of the data:

TOTEM Collab., EPL 101 (2013) 21002

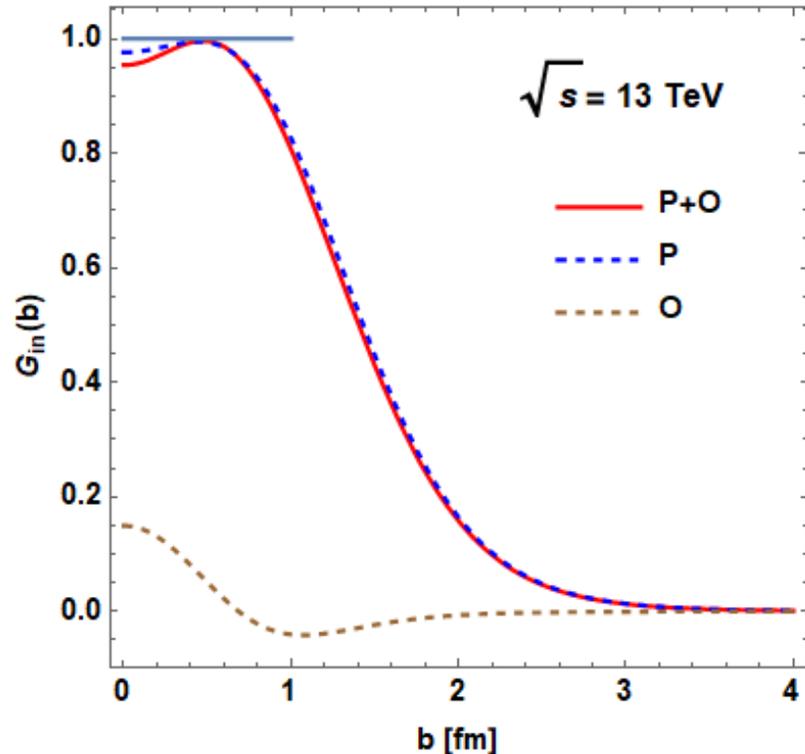
Fit of $d\sigma(t)/dt$ @ 2.76 TeV



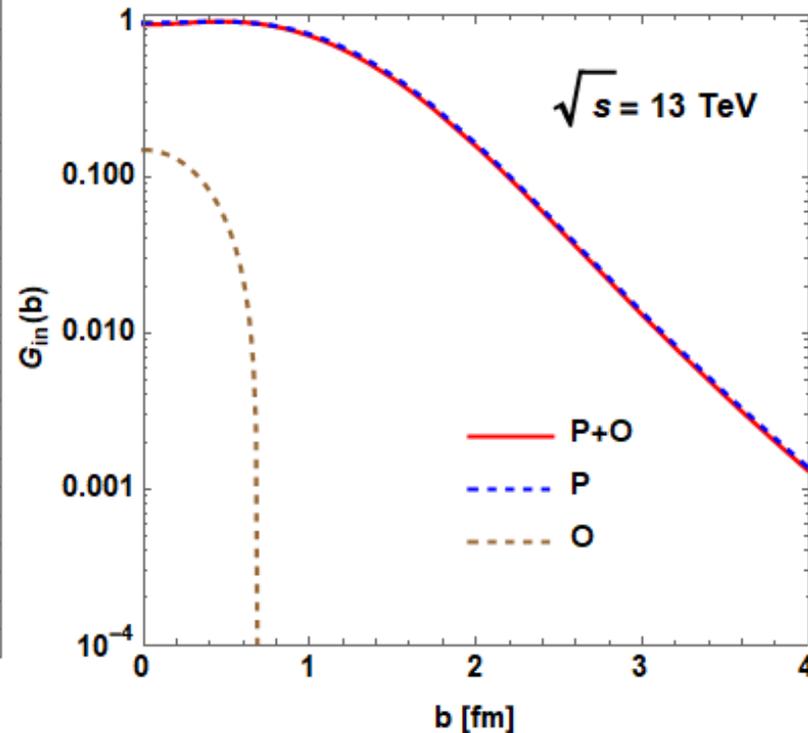
Normalized form of the differential cross section showing the low- $|t|$ "break"

Source of the data:
TOTEM Collab., CERN-EP-2018-341

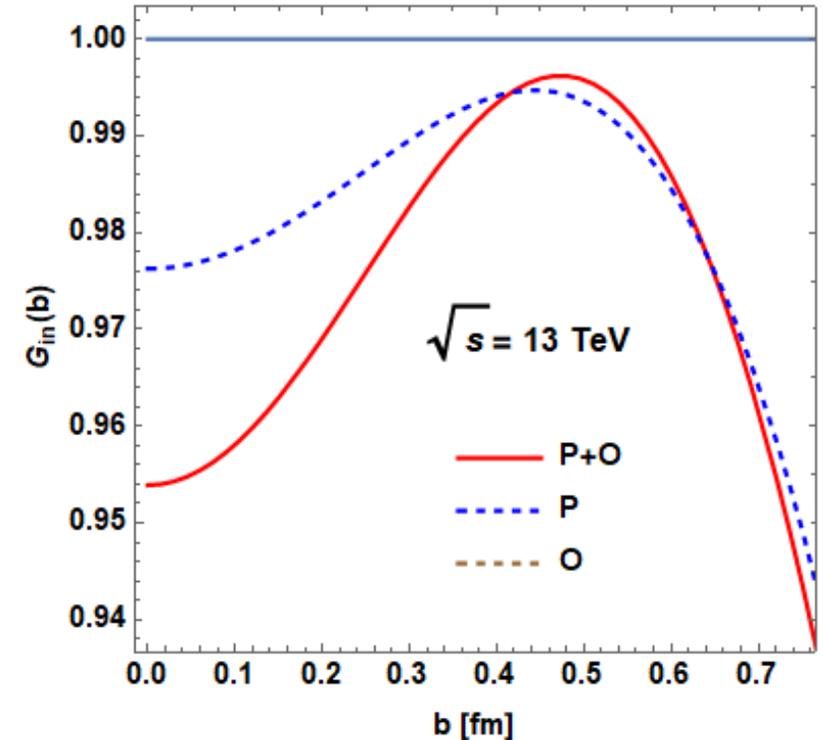
Inelastic overlap function @ 13 TeV



Calculated $G_{in}(b)$.



$G_{in}(b)$ with logarithmic vertical axis.

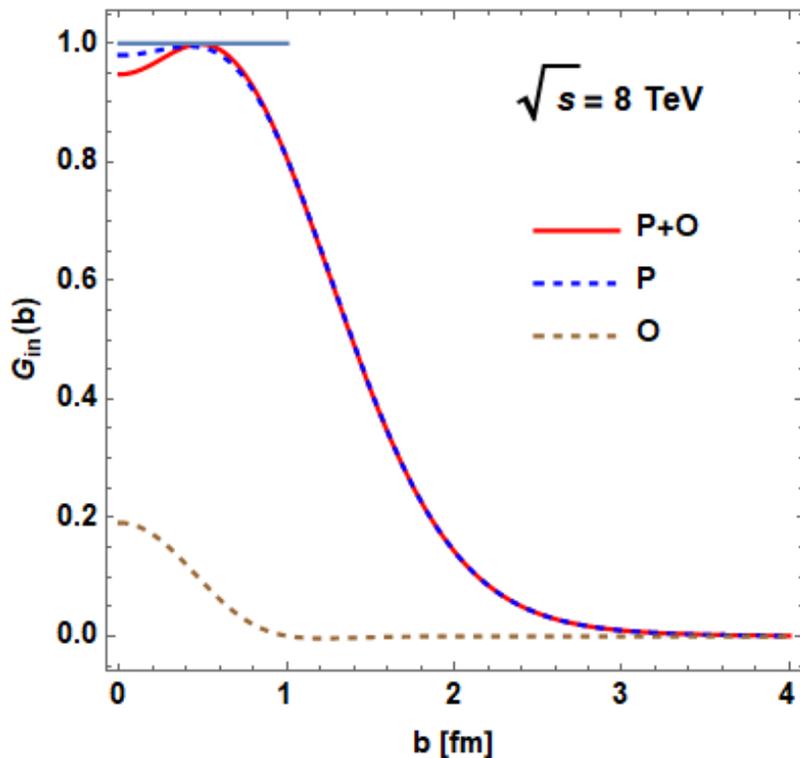


$G_{in}(b)$ enlarged for low b values.

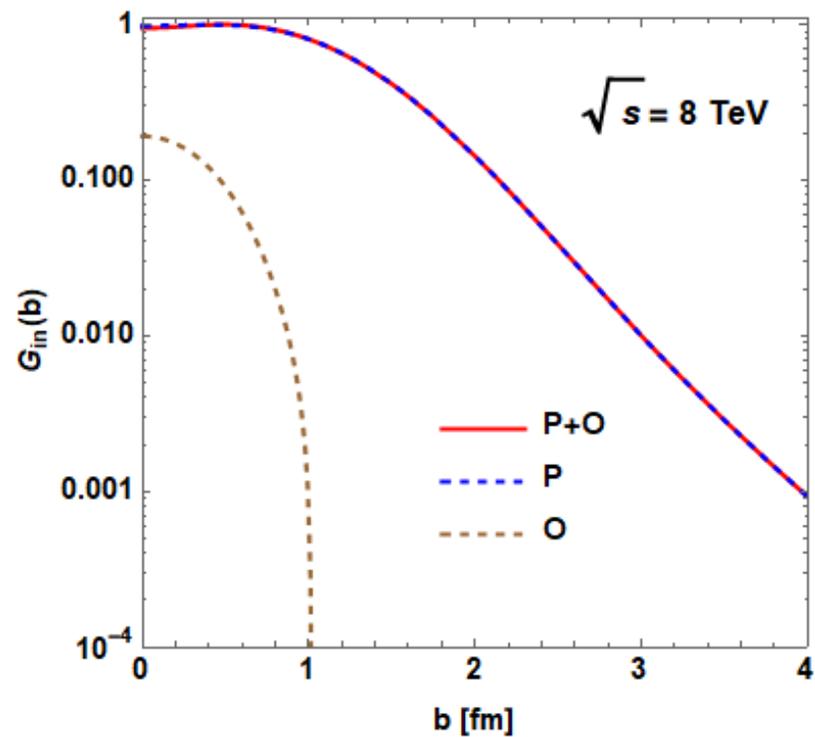
Similar results concerning the low- b G_{in} :

- I.M. Dremin, Physics 1 (2019) 33-39.
- T. Csörgő, R. Pasechnik, A. Ster, arXiv:1902.00109.
- W. Broniowski, L. Jenkovszky, E. Ruiz Arriola, I. Szanyi, Phys. Rev. D 98, 074012 (2018).

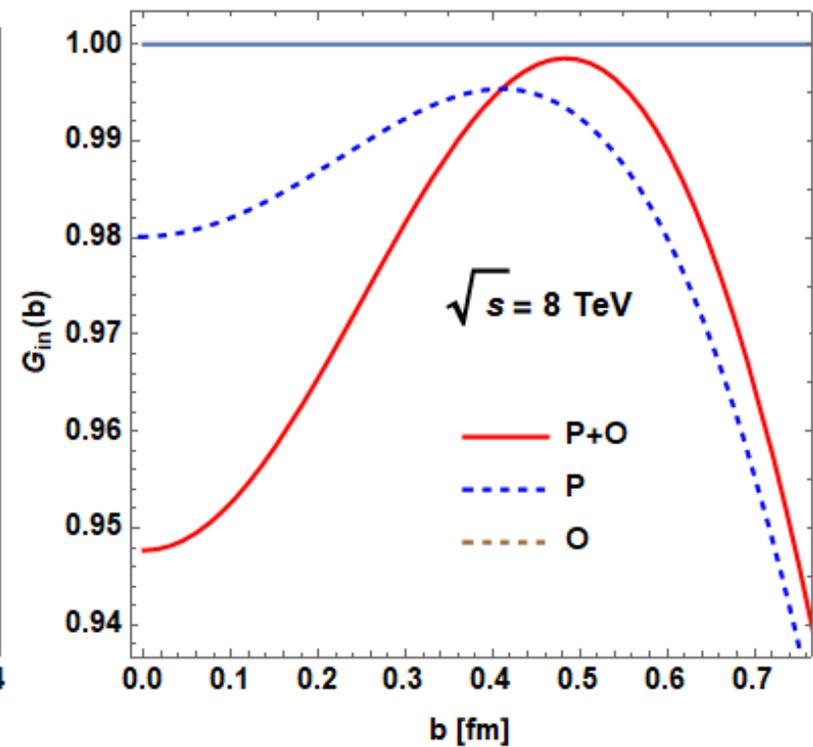
Inelastic overlap function @ 8 TeV



Calculated $G_{in}(b)$.

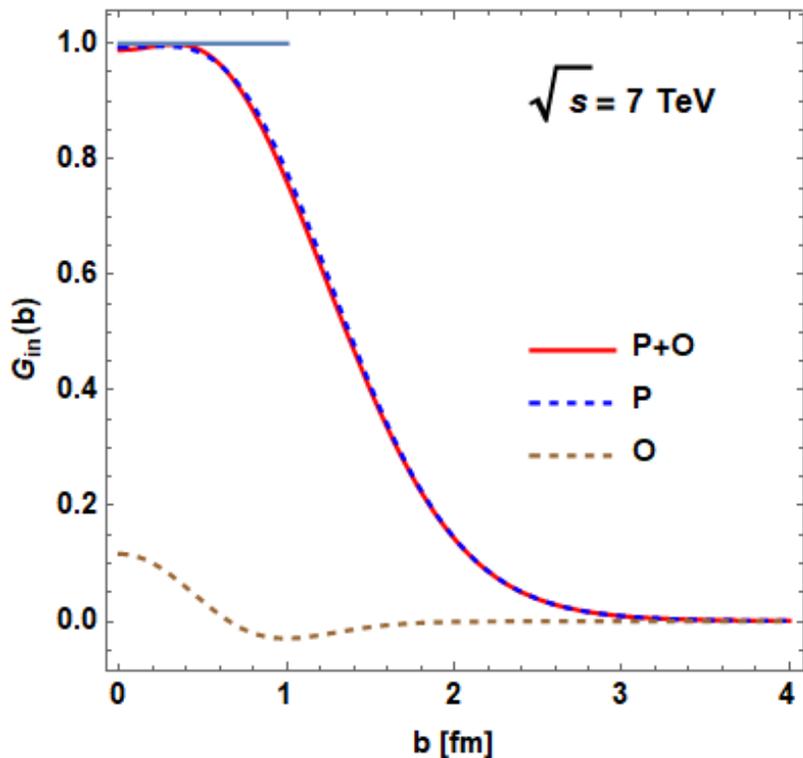


$G_{in}(b)$ with logarithmic vertical axis.

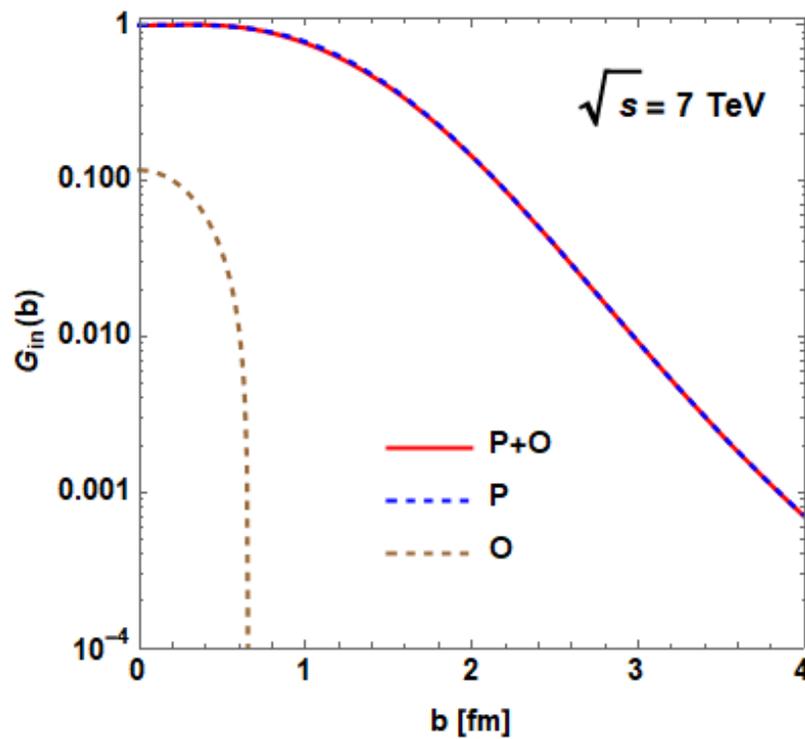


$G_{in}(b)$ enlarged for low b values.

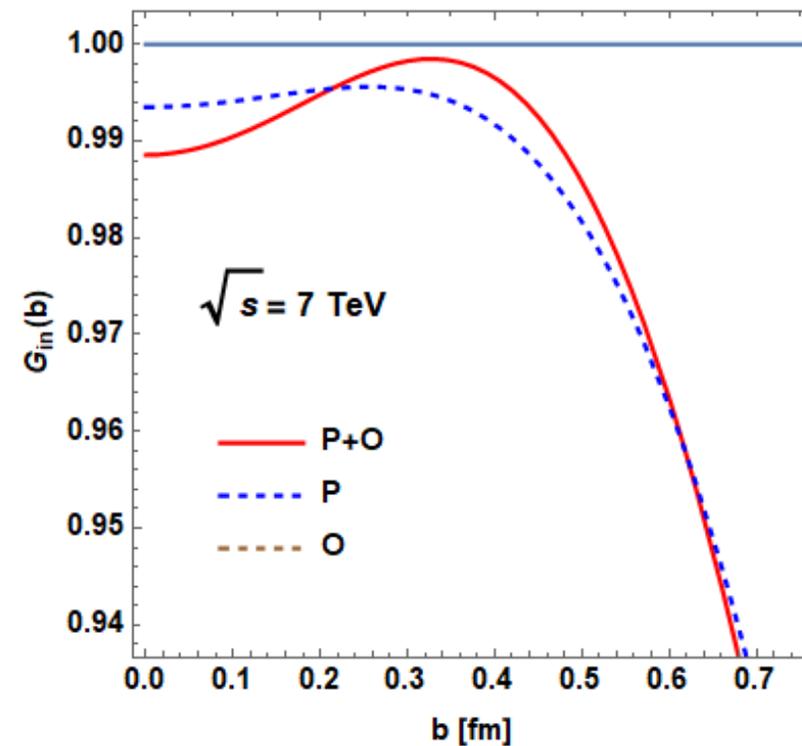
Inelastic overlap function @ 7 TeV



Calculated $G_{in}(b)$.

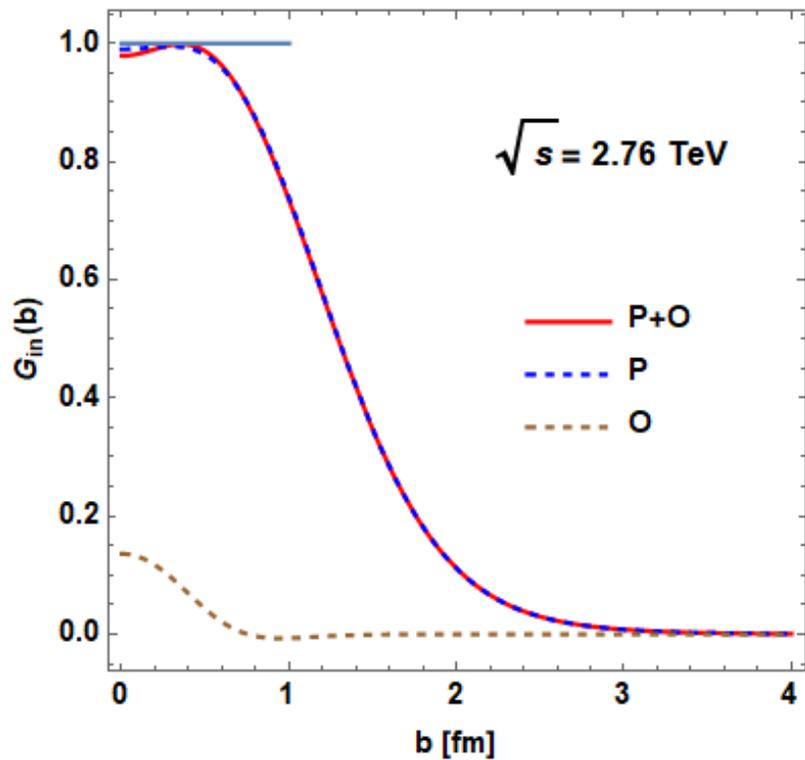


$G_{in}(b)$ with logarithmic vertical axis.

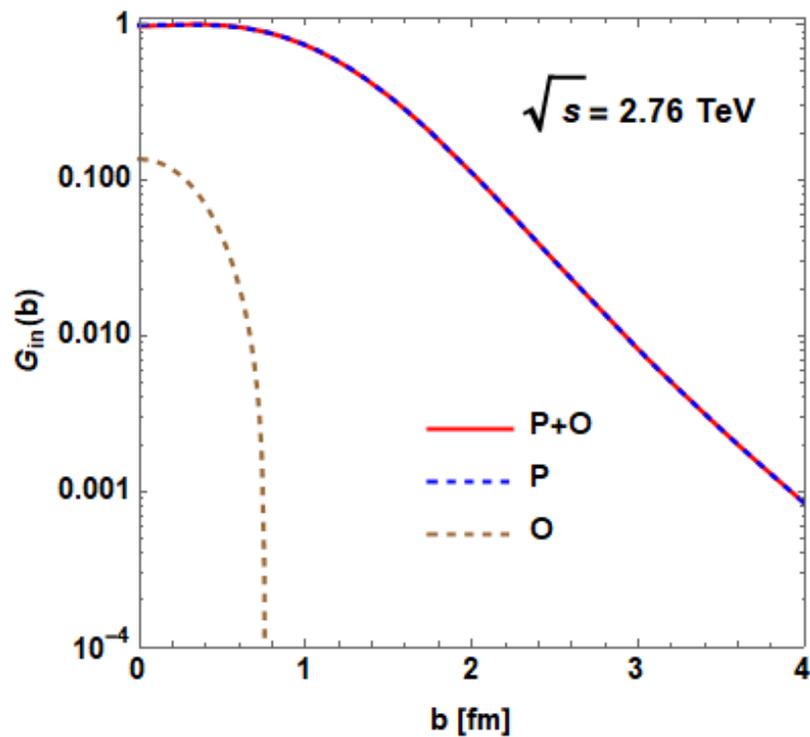


$G_{in}(b)$ enlarged for low b values.

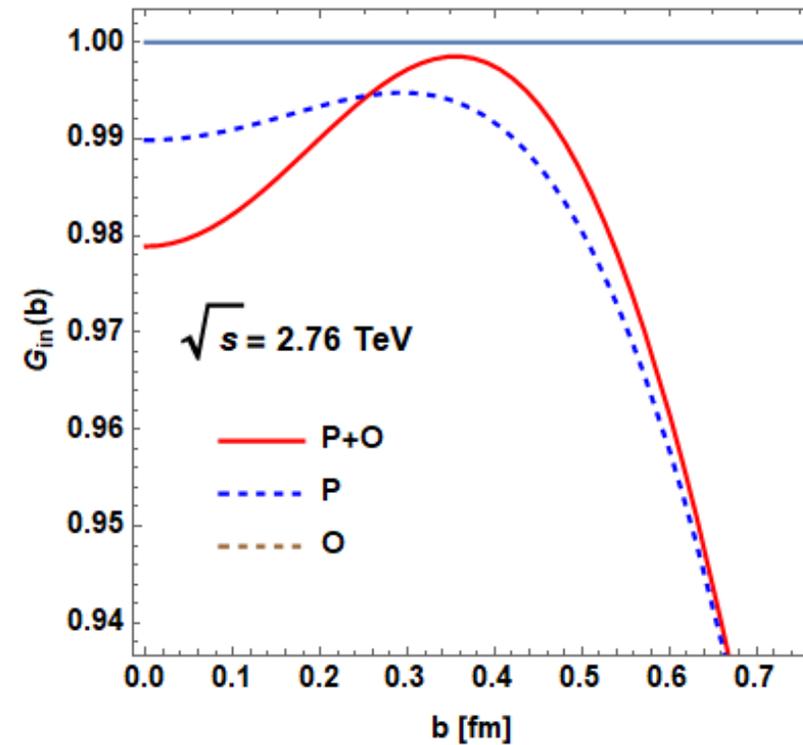
Inelastic overlap function @ 2.76 TeV



Calculated $G_{in}(b)$.

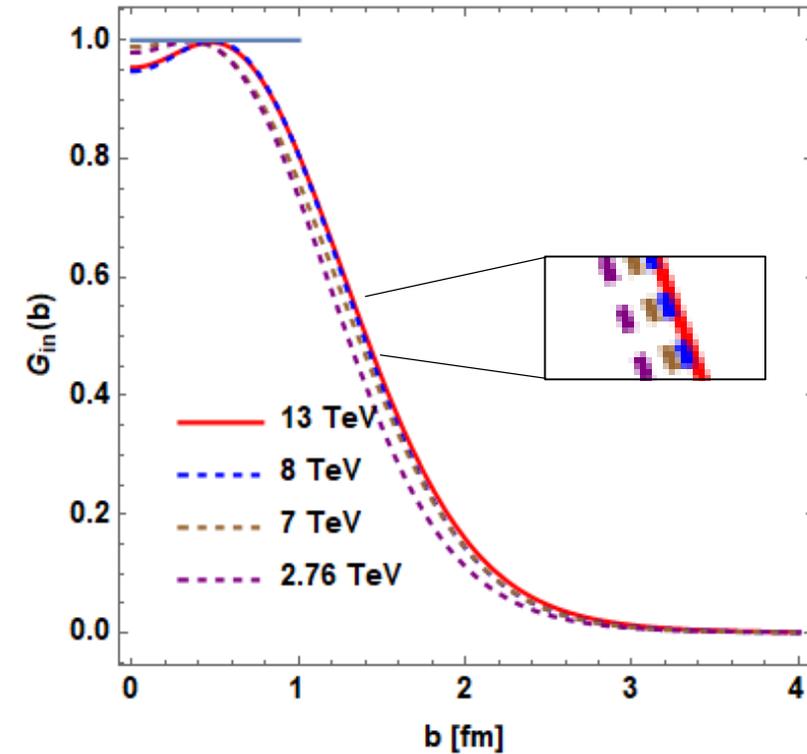


$G_{in}(b)$ with logarithmic vertical axis.

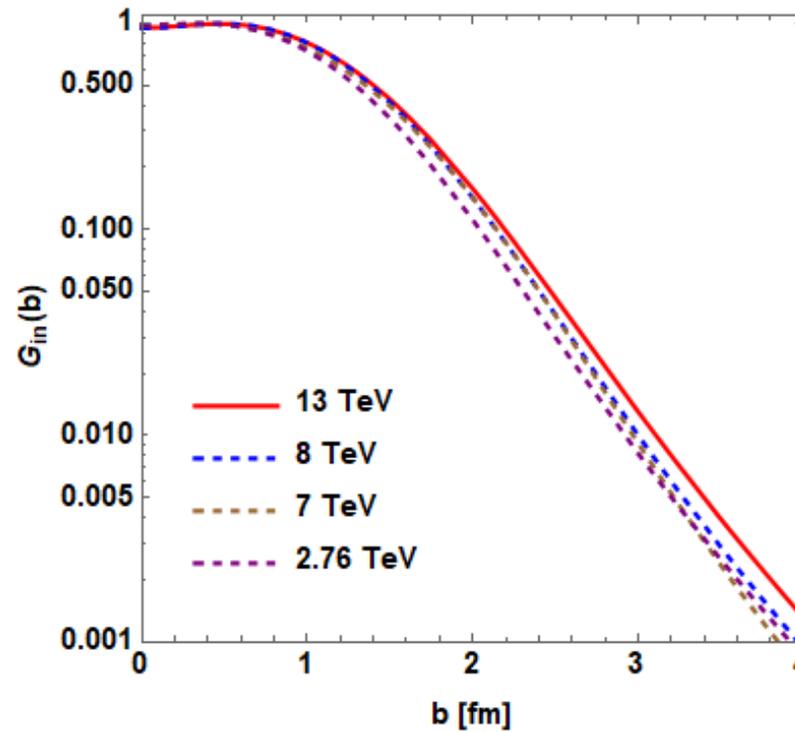


$G_{in}(b)$ enlarged for low b values.

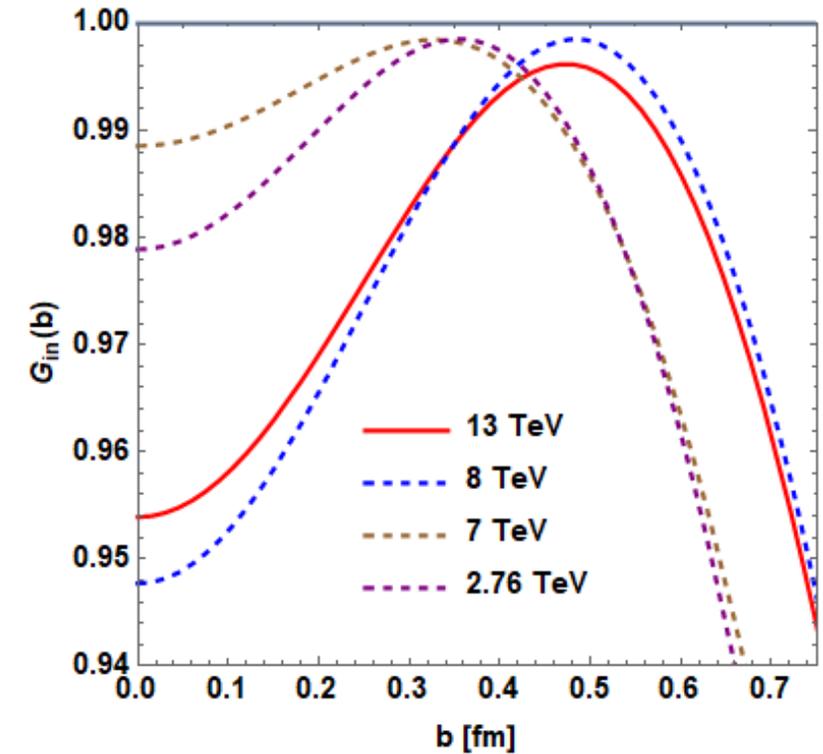
Comparison of $G_{in}(b)$ @ different energies



$G_{in}(b)$ calculated at different energies.



$G_{in}(b)$ calculated at different energies and illustrated with logarithmic vertical axis.



$G_{in}(b)$ calculated at different energies and enlarged for small b values.

Summary and conclusions

- Fits for the newest TOTEM proton-proton differential cross section data using a Regge model with dipole pomeron and odderon.
- Determination of the impact parameter amplitude.
- Calculation of inelastic overlap functions.
- Conclusions for the investigated energy range in the framework of the used model:
 - the proton is surrounded by pion cloud and its effective size is growing with energy
 - the interaction region of the colliding protons has a toroid-like shape (hollowness)
 - this shape is dominantly determined by the pomeron component of the amplitude to which the odderon gives a smaller contribution
 - the energy dependence of the minimum and the maximum of the G_{in} does not show a regularity
 - problems may arise from the fact that the exact t -dependent phase cannot be recovered from the experimental data

Thank you for your attention!

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