Impact of asymmetric bosonic dark matter on neutron star properties

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Bosonic DM in NS

GW170817: the first NS-NS merger

• Typical inspiral pattern

$$\begin{cases} M_{\rm tot} = 2.84^{+0.47}_{-0.09} M_{\odot} \\ \mathcal{M} = 1.188^{+0.04}_{-0.02} M_{\odot} \end{cases}$$

- Kilonova-associated
- Tidal deformability parameter extraction;



Figure 1: GW170817 and its γ counterpart signals. [LIGO/Virgo Collab 2017]]

(1)

Tidal disruption: a new constraints

• Tidal deformability parameter can be written as:

$$\Lambda = \frac{2}{3}k_2 \left(\frac{R_{\rm out}}{GM}\right)^5 \tag{2}$$

being $k_2 = k_2(\beta, R)$ the Love's number, R_{out} the outer-most radius and M the total mass.

- Masses extracted by the GW signal;
- Constraints on Λ, R as:

$$\begin{cases} \Lambda(1.4M_{\odot}) \le 800\\ R(1.4M_{\odot}) = 11.89^{+2.7}_{-3.3} \text{km} \end{cases}$$
(3)



Figure 2: Last time-steps before the merger phase.

[Radice et al., ApJ 869:130 (2018)]

An insight on the NS interior



Figure 3: (*Left*) Λ -Constraints on the two NSs of GW170817, (*Right*) Constraints derived from the same event on the Equation of State.

LIGO-Virgo Collaboration Phys. Rev. X 9, 011001 (2019)

Dark matter distribution

- Light dark matter (DM) particles $(m_{\chi} \simeq 100 \text{ MeV})$ tend to create diluted halo structures inside the Neutron Star (NS);
- Heavier DM particles $(m_{\chi} \gtrsim 1 \text{ GeV})$ form compact cores.



dark halo around a NS

Figure 4: Different configurations of DM structures in compact objects.

V. Sagun et al. arXiv:2111.13289 [astro-ph.HE]

TOV equations

2 TOV equations:

$$\frac{dp_B}{dr} = -\frac{(\epsilon_B + p_B)(M + 4\pi r^3 p)}{r^2 (1 - 2M/r)}$$

$$\frac{dp_D}{dr} = -\frac{(\epsilon_D + p_D)(M + 4\pi r^3 p)}{r^2 (1 - 2M/r)}$$

BM and DM are coupled only through gravity, and their energy-momentum tensors are conserved separately

total pressure $p(r) = p_B(r) + p_D(r)$

gravitational mass $M(r) = M_B(r) + M_D(r)$, where $M_j(r) = 4\pi \int_0^r \epsilon_j(r') r'^2 dr'$ (j=B,D)

Fraction of DM inside the star:

$$f_{\chi} = \frac{M_D(R_D)}{M_T}$$

 $M_T = M_B(R_B) + M_D(R_D)$ - total gravitational mass

O. Ivanytskyi, et al. PRD 102, 063028 (2020)

Effect of dark matter

• DM effects mass, radius and, therefore, tidal deformability of NSs

Tidal deformability parameter:

Mass-Radius diagram

$$\Lambda = \frac{2}{3}k_2 \left(\frac{R_{\rm out}}{GM}\right)^5 \qquad (4)$$



An example of an impact of fermionic DM on M-R relation and tidal deformability of NSs

V. Sagun, et al. arXiv:2111.13289 [astro-ph.HE]

Tidal deformabilities

Bosonic stars

- Considering bosonic DM particles, there's no Fermi pressure and the uncertainty principle is the only one preventing the collapse;
- Once captured, DM particles slow down and tend to thermalize in the core

$$r_{\rm th} = \left(\frac{9kT_c}{8\pi\rho_c m}\right)^{\frac{1}{2}} \tag{5}$$

$$t_{\rm th} = 0.2 {\rm yr} \left(\frac{m}{{\rm TeV}}\right)^2 \left(\frac{\sigma}{10^{-43} {\rm cm}^2}\right)^{-1} \left(\frac{T}{10^5 {\rm K}}\right)^{-1}$$
(6)

• The bosonic component can become self-gravitating

$$M_{\rm sg} > \frac{4\pi}{3} \rho_c r_{\rm th}^3 \tag{7}$$

or form a Bose-Einstein Condensate (BEC).

• When $r_{\rm BEC} \lesssim r_{\rm sch}$, a black hole is formed in the inner core.

C. Kouvaris, Adv.High Energy Phys. 2013, 856196 (2013)

Asymmetric Bosonic Dark Matter

The Lagrangian used includes **self-interacting** DM particles as:

$$\mathcal{L} = \frac{1}{2} (D_{\mu} \varphi)^* D^{\mu} \varphi - \frac{1}{2} m_{\varphi}^2 \varphi^* \varphi - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} - \frac{1}{2} m_{\omega}^2 \omega_{\mu} \omega^{\mu}$$
(8)

with $D_{\mu}\varphi = (\partial_{\mu} - ig\omega_{\mu})\varphi$.

Using a mean field approximation for ω , we recollect after (not so short) calculations an EoS of self-interacting Bose-Einstein condensate:

$$n = \frac{m_I^2}{2} \frac{\mu_{\varphi}^2 - m_{\varphi}^2}{\sqrt{2m_{\varphi}^2 - \mu_{\varphi}^2}}$$
(9)
$$p = \left(\frac{m_I}{2}\right)^2 \left[m_{\varphi}^2 - \mu_{\varphi}^2\sqrt{2m_{\varphi}^2 - \mu_{\varphi}^2}\right]$$
(10)

Chemical potential is limited $\mu_{\varphi} \in [m_{\varphi}, \sqrt{2}m_{\varphi}], \quad m_{\varphi}$ - boson mass $m_I = \frac{m_{\omega}}{q}$ - self-interaction scale

Induced Surface Tension (IST) EoS

$$\begin{cases} p = \sum_{i}^{all \ particles} [p_{id}(T, \mu_i - pV_i - \Sigma S_i + U_{at} \pm U_{sym}) + p_{id}(\mu_e) - p_{at} + p_{sym}] \\ \Sigma = \sum_{i}^{all \ particles} p_{id}(T, \mu_i - pV_i - \alpha \Sigma S_i + U_0)R_i \end{cases}$$

 p_{id} – pressure of the ideal gas for quantum statistics V – excluded volume, Σ – induced surface tension U_0, α – model parameters

• Thermodynamic consistency of the model:

$$\frac{\partial p_{int}}{\partial n_{id}} = n_{id} \tag{11}$$

• Parametrization of the mean field potential:

$$U_{at} = -C_d^2 n_{id}^{\kappa} \tag{12}$$

• Symmetry energy pressure:

$$P_{sym}(n) = \frac{A_{sym}n^2}{[1 + (B_{sym}n)^2]^2}$$
(13)

VVS, et al., Nucl. Phys. A, 924, 24 (2014)

A. Ivanytskyi et al., PRC 97, 064905 (2018)

Bosonic DM in NS

Gravitational instability of DM-admixed NS



- consider fixed central densities n_D^c and n_B^c
- solve TOV $\Rightarrow R_D$ and R_B
- vary n_D^c to get new R_D and R_B
- along the curves central density of DM component is varying

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Schwarzschild radius is defined as:

$$R_{\rm Sch} = 2GM \tag{14}$$

To avoid an instability against complete gravitational collapse the star's radius should be of the order of the last stable Kepler orbit $R_{SO} = 3R_{Sch}$ around a BH of Schwarzschild radius.

The radius of the BEC equals to R_D

E. Giangrandi et al., in Prep., (2022)



- We considered a novel model of self-interacting bosonic DM and tested an effect of particle's mass, coupling constant and DM fraction of compact star properties
- We analyzed conditions of gravitational instability in the core of a star that appears at R_D equal to the last stable Kepler orbit
- We admit a possibility of a DM condensate driven BH formation inside a compact star.

Thank you for attention!